Ergodic Theorems: Suppose

$$
\mathcal{L}=\frac{1}{2} \sum_{i, j} a_{i, j}(x) D_{i, j}+\sum_{j} b_{j}(x) D_{j}
$$

is the generator of a process. Assume that $\left\{a_{i, j}\right\}$ are continuous and strictly positive definite, i.e. $\left\{a_{i, j}\right\}$ is positive definite for each $x \in R^{d}$. An invariant probability measure is one for which

$$
\int p(t, x, A) \mu(d x)=\mu(A)
$$

for every Borel set $A \subset R^{d}$. First $\mu$ must be absolutely continuous, i.e. $\mu(d y)=\phi(y) d y$ for some $\phi \in L_{1}\left(R^{d}\right)$. The proof depends on the fact that $p(t, x, d y)$ in fact has a density $p(t, x, y) d y$ for every $t>0$, and $x \in R^{d}$. We would like to prove the following.

1. Uniqueness. $\mu(d x)=\phi(x) d x$ is unique. $\phi(x)>0$ a.e.
2. Law of Large Numbers. With probability 1 with respect to any $P_{x}$,

$$
\lim _{t \rightarrow \infty} \frac{1}{t} \int_{0}^{t} f(x(s)) d s=\int f(y) \phi(y) d y
$$

for any bounded measurable function $f(\cdot)$.

## 3. Convergence to equilibrium.

$$
\lim _{t \rightarrow \infty} \int|p(t, x, y)-\phi(y)| d y=0
$$

4. Central Limit Theorems. If $\int f(y) \phi(y) d y=0$, under additional assumptions

$$
\frac{1}{\sqrt{t}} \int_{0}^{t} f(x(s)) d s
$$

has an asymptotic normal distribution for large $t$.
5. Large Deviations from the Ergodic Theorem. Study the large deviation behavior

$$
P_{x}\left[\frac{1}{t} \int_{0}^{t} f(x(s)) d s \in A\right]
$$

Uniqueness: Uniqueness and ergodicity are related. Let $\mu$ be an invriant probability measure for the Markov process. This means that if we start with initial distribution $\mu$ the $P_{\mu}=\int P_{x} d \mu(x)$ is staionary process for $t \geq 0$ and can be easily extended as a stationary process for $-\infty<t<\infty$. If $P_{\mu}$ is not ergodic with respect to trnaslations, then there exists a trnaslation invariant set $E$ on the space $C\left[(-\infty, \infty) ; R^{d}\right]$ of trajectories such that $0<P_{\mu}(A)<1$. We will show now that $E$ can be chosen to be of the form $x(0) \in A$ for some Borel set. Note that any measurable set $E$ can be approximated by sets measurable with
respect to the $\sigma$-field $\sigma\{x(t):|t| \leq T\}$ and if it is translation invariant can be approximated by a set from $\sigma\{x(t): T \leq t \leq 2 T\}$ as well as from $\sigma\{x(t):-2 T \leq t \leq-T\}$. In particular it is in the past $\sigma\{x(t): t \leq-T\}$ as well as the future $\sigma\{x(t): t \geq T\}$. But for a Markov process given the present $x(0)$ the past and the future are independent. This means given $x(0)$ the set $E$ is indpendent of itself. Therefore $P[E \mid x(0)]=0$ or 1 . Now $E$ can be identified with the set $\{x(\cdot): P[E \mid x(0)]=1\}$. If there are two invariant measures $\mu_{1}, \mu_{2}$, then clearly $P_{\mu}$ with $\mu=\frac{1}{2}\left[\mu_{1}+\mu_{2}\right]$ can not be ergodic.

In particular if the invariant measure $\mu$ is not unique we can find a partition of $R^{d}=A \cup A^{c}$ such that the restrictions of $\mu$ to $A$ and $A^{c}$ are both invariant. This means for every $t>0, p(t, x, A)=1$ for a.e. $x \in A$ and $p\left(t, x, A^{c}\right)=1$ for a.e $x \in A^{c}$.

We note that, from PDE , we can obtain that $p(t, x, A)$ is a jointly continuous function of $t, x$ for $t>0$. In partcular $y: p(t, y, A)>\frac{1}{2}$ is an open set $U$, and for almost all $x \in A^{c}$

$$
p(2 t, x, A)=\int p(t, x, d y) p(t, y, A) \geq \frac{1}{2} p(t, x, U)=0
$$

Therefore to prove uniqueness it suffices to show that given $x \in R^{d}$ and any open ball $B$, $p(t, x, B)>0$ for some $t>0, x \in R^{d}$ any ball $B$. Suppose $P_{x}[x(t) \in B]=0$ for all $t$. By Girsanov's theorem the same is true if we change the drift, and also if we rescale time. Therefore

$$
Q_{x}^{\epsilon}[x(t) \in B]=0
$$

where $Q_{x}^{\epsilon}$ corresponds to

$$
\mathcal{L}_{\epsilon}=\frac{\epsilon}{2} \sum_{i, j} a_{i, j}(x) D_{i, j}+\sum_{j}\left[\epsilon b_{j}(x)+c_{j}(x)\right] D_{j}
$$

with $\left\{c_{j}(x)\right\}$ having compact support. If we let $\epsilon \rightarrow 0$ the same is true for the limit. But the limit is just an O.D.E and we can pick our coeffiicients!. This can not be true.

This proves the uniqueness of $\mu$ and the ergodicity of $P_{\mu}$. It is not hard to see that from

$$
\mu(A)=\int p(t, x, A) d \mu(x)
$$

and the existence of a density $p(t, x, y)$ we can conclude that $\mu$ in fact has a density $\pi(x)$. The positivity of $\phi(x)$ again come from PDE. If $\int_{A} \phi(x) d x=0$ for a set with positive Lebesgue meausre, then $p(t, x, A)$ would be 0 for all $x$ and $t>0$, and the PDE estimate shows the convergence of $p(t, x, A) \rightarrow \chi_{A}(x)$ in $L_{p}$ as $t \rightarrow 0$. Another way of verifying uniqueness is by a direct argument. Let $\pi(x, y)=p(1, x, y)$. If $\mu$ is invariant then, it is invariant for $\pi(x, y)$ as well. Hence

$$
\int \pi(x, y) d y \mu(d x)=\mu(d y)
$$

This implies that $\mu(d x)=\phi(x) d x$ and

$$
\int \pi(x, y) \phi(x) d x=\phi(y)
$$

If there are two solutions let us denote their difference by $\psi(x)$. Then $\int \psi(x) d x=0$,

$$
\int \psi(x) \pi(x, y) d x=\psi(y)
$$

and

$$
\int|\psi(y) d y|=\int\left|\int \psi(x) \pi(x, y) d x\right| d y \leq \iint|\psi(x)| \pi(x, y) d x d y=\int|\psi(x)| d x
$$

In particular for almost all $y$,

$$
\left|\int \psi(x) \pi(x, y) d x\right|=\int|\psi(x)| \pi(x, y) d x
$$

This can not happen unless $\psi(x)$ is of one sign. But it has integral equal to 0 . Hence $\psi=0$.

Ergodic Theorem. Because of ergodicity we know that with probability 1 with respect to any $P_{\mu}$,

$$
\lim _{t \rightarrow \infty} \frac{1}{t} \int_{0}^{t} f(x(s)) d s=\int f(y) \phi(y) d y
$$

for any bounded measurable function $f(\cdot)$. On the other hand form the positivity of $p(t, x, y)$, for $t>0, p(t, x, y)$ and $\mu$ are mutually absolutely continuous. By the Markov property if we ignore a small initial segment of time then the almost sure convergence w.r.t. $P_{\mu}$ gives a.e. convergence w.r.t $P_{x}$ for any $x$.

Convergence to equilibrium. This is not much different from uniqueness. Let us denote by

$$
\delta(t, x)=\int|p(t, x, y)-\phi(y)| d y
$$

then

$$
\begin{aligned}
\delta(t+s, x) & =\int\left|\int\left[p\left(t, x, z^{\prime}\right)-\phi\left(z^{\prime}\right)\right] p\left(s, z^{\prime}, z\right) d z^{\prime}\right| d z \\
& \leq \int\left|p\left(t, x, z^{\prime}\right)-\phi\left(z^{\prime}\right)\right| p\left(s, z^{\prime}, z\right) d z d z^{\prime} \\
& =\delta(t, x, y)
\end{aligned}
$$

proving $\delta(t, x)$ is $\downarrow$ in $t$ for fixed $x$. It is therefore sufficient to prove that $\delta(n, x) \rightarrow 0$ as $n \rightarrow \infty$. We have a Markov Chain with almost surely positive transition density $\pi(x, y)=p(1, x, y)$ satisfying, $x \rightarrow \pi(x, \cdot)$ is continuous as a map into $L_{1}\left(R^{d}\right)$ that admits an invariant probability density $\phi(x)$. We need to prove that for compact $K$,

$$
\lim _{n \rightarrow \infty} \sup _{x \in K} \int\left|\pi^{n}(x, y)-\phi(y)\right| d y
$$

The basic idea here is coupling. Let us take two copies with transition probability

$$
\begin{aligned}
\widehat{\pi}\left(x_{1}, x_{2} ; d y_{1}, d y_{2}\right)= & \int \min \left\{\pi\left(x_{1}, y\right) \pi\left(x_{2}, y\right)\right\} \delta\left(y-y_{1}\right) \times \delta\left(y-y_{2}\right) d y \\
+ & \int\left[\pi\left(x_{1}, y\right)-\min \left\{\pi\left(x_{1}, y\right) \pi\left(x_{2}, y\right)\right\}\right] \delta\left(y-y_{1}\right) d y \\
& \times \int\left[\pi\left(x_{2}, y\right) \min \left\{\pi\left(x_{1}, y\right) \pi\left(x_{2}, y\right)\right\}\right] \delta\left(y-y_{2}\right) d y
\end{aligned}
$$

In words, if $g(y)=\min \min \left\{\pi\left(x_{1}, y\right) \pi\left(x_{2}, y\right)\right\}$ the two components jump on to the diagonal $y=y_{1}=y_{2}$ with density $g(y) d y$ and jump indpendently with the remaining probability. Once they are on the diagonal they stay on the diagonal. It is not hard to see that each component evolves according to $\pi$. Therefore

$$
\int\left|\pi^{n}\left(x_{1}, y\right)-\pi^{n}\left(x_{2}, y\right)\right| d y \leq \widehat{\pi}^{n}\left(x_{1}, x_{2} ;\left\{y_{1} \neq y_{2}\right\}\right)=\delta_{n}\left(x_{1}, x_{2}\right)
$$

and if $r_{n}\left(x_{1}, x_{2}, d y_{1}, d y_{2}\right)$ is the part of $\widehat{\pi}^{n}$ away from the diagonal

$$
\delta_{n+1}\left(x_{1}, x_{2}\right)=\delta_{n}\left(x_{1}, x_{2}\right)-2 \int \min \left\{\pi\left(y_{1}, y\right), \pi\left(y_{2}, y\right)\right\} r_{n}\left(x_{1}, x_{2}, d y_{1}, d y_{2}\right) d y
$$

Since $\delta_{n} \geq 0$ we must have

$$
\int \min \left\{\pi\left(y_{1}, y\right), \pi\left(y_{2}, y\right)\right\} r_{n}\left(x_{1}, x_{2}, d y_{1}, d y_{2}\right) d y \rightarrow 0
$$

as $n \rightarrow \infty$. Since $\pi(x, y)>0$ a.e, this is possible only if

$$
\int_{\substack{\left|y_{1}\right| \leq \ell \\\left|y_{2}\right| \leq \ell}} r_{n}\left(x_{1}, x_{2}, d y_{1}, d y_{2}\right) \rightarrow 0
$$

On the other hand if the total mass of $r_{n}$ does not go to 0 , we must have

$$
\lim _{\ell \rightarrow \infty} \liminf _{n \rightarrow \infty} r_{n}\left(x_{1}, x_{2},\left\{\left|y_{1}\right| \geq \ell \cup\left|y_{2}\right| \geq \ell\right\}\right) \geq c>0
$$

Since

$$
r_{n}\left(x_{1}, x_{2},\left\{\left|y_{1}\right| \geq \ell \cup\left|y_{2}\right| \geq \ell\right\}\right) \leq \pi^{n}\left(x_{1},\left|y_{1}\right| \geq \ell\right)+\pi^{n}\left(x_{2},\left|y_{2}\right| \geq \ell\right)
$$

this contradicts the ergodic theorem

$$
\frac{1}{n} \sum_{j=1}^{n} \pi^{j}(x, A) \rightarrow \mu(A)
$$

Central Limit Theorem. Consider

$$
A(t)=\int_{0}^{t} f(x(s) d s
$$

If $\int f(x) d \mu(x)=0$, one can expect to prove the C.L.T that states

$$
\frac{A(t)}{\sqrt{t}} \simeq N\left(0, \sigma^{2}\right)
$$

The proof depends on ability to solve

$$
\mathcal{L} u=-f
$$

with a nice $u$. This is the Fredholm alternative. Then by Itô's formula

$$
u(x(t))-u(x)=\int_{0}^{t} \sigma(x(s))(\nabla u)(x(s)) d \beta(s)-\int_{0}^{t} f(x(s)) d s
$$

Therefore

$$
A(t)=\int_{0}^{t} \sigma(x(s))(\nabla u)(x(s)) d \beta(s)-u(x(t))-u(x)
$$

Dividing by $\sqrt{t}$ if we assume that $\frac{u(x(t))}{\sqrt{t}} \rightarrow 0$ as $t \rightarrow \infty$, we need only analyse

$$
\frac{1}{\sqrt{t}} \int_{0}^{t} \sigma(x(s))(\nabla u)(x(s)) d \beta(s)=\widehat{\beta}\left(\frac{1}{t} \int_{0}^{t}<a(x(s)) \nabla u(x(s)), \nabla u(x(s))>d s\right)
$$

and by the ergodic theorem this converges to $\widehat{\beta}\left(\sigma^{2}\right) \simeq N\left(0, \sigma^{2}\right)$ with

$$
\sigma^{2}=\int<a(x)(\nabla u)(x),(\nabla u)(x)>d \mu(x)
$$

One can work with approximate solutions of $\mathcal{L} u=-f$ namely

$$
\lambda u_{\lambda}-\mathcal{L} u_{\lambda}=f
$$

and reduce the CLT to the behavior of $u_{\lambda}$ as $\lambda \rightarrow 0$. One needs

$$
\lambda \int\left\|u_{\lambda}\right\|^{2} d \mu \rightarrow 0
$$

and

$$
\left.\int<a(x)\left(\nabla u_{\lambda}(x)-g(x)\right),\left(\nabla u_{\lambda}(x)-g(x)\right)>d \mu_{( } x\right) \rightarrow 0
$$

for some $g$. One then applies Itô's formula to $u_{\lambda}(x)$ and chooses $\lambda=t^{-1}$. The errors are easily controlled.

To see this we need to control

$$
X_{\lambda}(t)=\frac{1}{\sqrt{t}} \int_{0}^{t} \lambda u_{\lambda}(x(s)) d s
$$

Starting from the equlibrium $\mu$

$$
\left\|X_{\lambda}(t)\right\|_{2} \leq \lambda \sqrt{t}\left\|u_{\lambda}\right\|_{2} \leq \sqrt{\lambda}\left\|u_{\lambda}\right\|_{2} \rightarrow 0
$$

As for the stochastic integral part we have the CLT valid for

$$
\frac{M_{\lambda}(t)}{\sqrt{t}}=\frac{1}{\sqrt{t}} \int_{0}^{t} \sigma(x(s))\left(\nabla u_{\lambda}\right)(x(s)) d \beta(s)
$$

The difference between between $M_{\lambda_{1}}(t)$ and $M_{\lambda_{2}}(t)$ is easily estimated by

$$
\begin{aligned}
& \frac{1}{t} E\left[\left[M_{\lambda_{1}}(t)-M_{\lambda_{2}}(t)\right]^{2}\right] \\
& \quad=\int<a(x)\left(\nabla u_{\lambda_{1}}(x)-\nabla u_{\lambda_{2}}(x)\right),\left(\nabla u_{\lambda_{1}}(x)-\nabla u_{\lambda_{2}}(x)\right)>d \mu(x) \rightarrow 0
\end{aligned}
$$

A special situation is when $\mathcal{L}$ is self adjoint with respect to $\mu$. Then one defines the Dirichlet norm by

$$
\|u\|_{1}^{2}=\int<a(x)(\nabla u)(x),(\nabla u)(x)>d \mu(x)
$$

and the dual norm

$$
\|f\|_{-1}=\sup _{u} \frac{\int f(x) u(x) d \mu}{\|u\|_{1}}
$$

The finiteness of $\|f\|_{-1}$ implies both the conditions needed for CLT and this is easily seen by the spectral theorem. If $k(d \sigma)$ is the spectral measure of $f$ with respect to $\mathcal{L}$, then the assumption is

$$
\int_{0+}^{\infty} \frac{1}{\sigma} d k(\sigma)<\infty
$$

The conditions to be proved are

$$
\lim _{\lambda \rightarrow 0} \int_{0+}^{\infty} \frac{\lambda}{(\lambda+\sigma)^{2}} k(d \sigma)=0
$$

and

$$
\lim _{\substack{\lambda_{1} \rightarrow 0 \\ \lambda_{2} \rightarrow 0}} \int_{0+}^{\infty}\left|\frac{1}{\left(\lambda_{1}+\sigma\right)}-\frac{1}{\left(\lambda_{2}+\sigma\right)}\right|^{2} \sigma d k(\sigma)=0
$$

Both are consequences of the dominated convergence theorem.
Large Deviations. We wish to estimate

$$
P_{x}\left[\frac{1}{t} \int_{0}^{t} f(x(s)) d s \geq \ell\right]
$$

where $\ell>\int f(x) d \mu(x)$. Does this decay exponentially? If so what is the rate? We will do one example. Brownian motion on the circle. The operator is $\mathcal{L}=\frac{1}{2} \frac{d^{2}}{d x^{2}}$ acting on periodic functions of period 1. The invariant measure is just Lebesgue measure. $\phi(x) \equiv 1$.

One way to estimate the rate is to try and compute

$$
\lim _{t \rightarrow \infty} \frac{1}{t} \log E_{x}\left[\sigma \int_{0}^{t} f(x(s)) d s\right]=F(\sigma)
$$

Then by Tchebechev's inequality

$$
P_{x}\left[\frac{1}{t} \int_{0}^{t} f(x(s)) d s \geq \ell\right] \leq e^{-\sigma \ell t+t F(\sigma)+o(t)}
$$

and

$$
\lim \sup \frac{1}{t} \log P_{x}\left[\frac{1}{t} \int_{0}^{t} f(x(s)) d s \geq \ell\right] \leq F(\sigma)-\ell \sigma
$$

Optimizing over $\sigma>0$,

$$
\lim \sup \frac{1}{t} \log P_{x}\left[\frac{1}{t} \int_{0}^{t} f(x(s)) d s \geq \ell\right] \leq-\sup _{\sigma>0}[\ell \sigma-F(\sigma)]
$$

It is not hard to see that $F^{\prime}(0)=\int f(x) d x$, and therefore for $\ell>\int f(x) d x$,

$$
\sup _{\sigma>0}[\ell \sigma-F(\sigma)]=\sup _{\sigma}[\ell \sigma-F(\sigma)]
$$

The same is true in the other side. The large deviation rate function can be then

$$
I(\ell)=\sup _{\sigma}[\ell \sigma-F(\sigma)]
$$

What is $F(\sigma)$ ? We can look at Feynman-Kac formula and

$$
u(t, x)=E_{x}\left[\sigma \int_{0}^{t} f(x(s)) d s\right]
$$

solves

$$
u_{t}=\frac{1}{2} u_{x x}+\sigma f u ; u(0, x)=1
$$

If we denote the eigen values and eigen functions of $\frac{1}{2} u_{x x}+\sigma f u$ by $\left\{\lambda_{j}, \phi_{j}\right\}$, then

$$
u(t, x)=\sum_{j} \phi_{j}(x) e^{t \lambda_{j}} \int \phi_{j}(y) d y
$$

The eigenvalues tend to $-\infty$ and the one that counts will be the largest one, the principal eigen value. Note that $\phi_{1}(y)>0$. Hence

$$
F(\sigma)=\lambda_{1}(\sigma f)=\sup _{\|g\|_{2}=1}\left[\sigma \int f(x) g^{2}(x) d x-\frac{1}{2} \int\left|g_{x}\right|^{2} d x\right]
$$

We can rewrite this as

$$
F(\sigma)=\lambda_{1}(\sigma f)=\sup _{\substack{g \geq 0 \\ g(x) d x=1}}\left[\sigma \int f(x) g(x) d x-\frac{1}{8} \int \frac{\left|g_{x}\right|^{2}}{g} d x\right]
$$

This has an interpretation. Instead of $\frac{1}{2} \frac{d^{2}}{d x^{2}}$ modify with a drift

$$
\mathcal{L}_{g}=\frac{1}{2} \frac{d^{2}}{d x^{2}}+\frac{g_{x}}{2 g} \frac{d}{d x}
$$

This has invariant density $g$. Now the erigodic theorem will produce $g$ averages. The process then will behave like a process with invariant density $g$. Call this process $Q_{g, x}$.

$$
P_{x}[A]=\int_{A} \frac{d P_{x}}{d Q_{g, x}} d Q_{g, x}
$$

Our set $A$ will have $Q_{g, x}$ measure 1. We need to concentrate on how small the derivative $\frac{d P_{x}}{d Q_{g, x}}$ can be on $A$. By Girsanov formula

$$
\log \frac{d Q_{g, x}}{d P_{x}}=\int \frac{g_{x}}{2 g}(x(s)) d x(s)-\frac{1}{8} \int \frac{g_{x}^{2}(x(s))}{g^{2}(x(s))} d s
$$

The ergodic theorem under $Q_{g, x}$ will produce a limit

$$
\frac{1}{t} \log \frac{d Q_{g, x}}{d P_{x}} \rightarrow \frac{1}{4} \int \frac{g_{x}^{2}}{g} d x-\frac{1}{8} \int \frac{g_{x}^{2}}{g} d x=\frac{1}{8} \int \frac{g_{x}^{2}}{g} d x
$$

giving with the help of Jensen's inequality a lower bound of

$$
\exp \left[-\frac{t}{8} \int \frac{g_{x}^{2}}{g} d x\right]
$$

for a Brownian Motion to behave like a process with invariant density $g$ for a duration $t$. This provides for our original estimate

$$
I(\ell)=\inf _{\substack{g \geq 0, \int g d x=1 \\ \int f(x) g(x) d x=\ell}}\left[\frac{1}{8} \int \frac{g_{x}^{2}}{g} d x\right]
$$

