

Suppose u is a smooth function on R^d with compact support or decays fast enough and $\Delta u = f$, according to a theorem of Calderon and Zygmund, for $1 < p < \infty$ there is a constant $C(d, p)$ such that

$$\left\| \frac{\partial^2 u}{\partial x_i \partial x_j} \right\|_p \leq C(d, p) \|f\|_p$$

The heat kernel in d dimension is given by.

$$p(t, x, y) = \frac{1}{(2\pi)^{\frac{d}{2}}} \exp\left[-\frac{(x-y)^2}{2t}\right]$$

the semigroup by

$$(T_t f)(x) = \int_{R^d} f(y) p(t, x, y) dy$$

and R_λ the resolvent by

$$u(x) = (R_\lambda f)(x) = \int_0^\infty \int_{R^d} e^{-\lambda t} f(y) p(t, x, y) dy dt = \int_0^\infty e^{-\lambda t} (T_t f)(x) dt$$

It solves the equation

$$\lambda u - \frac{1}{2} \Delta u = f$$

Since $\|T_t f\| \leq \|f\|_p$ for every $t > 0$ and $1 \leq p \leq \infty$, $\|\lambda R_\lambda f\|_p \leq \|f\|_p$. $\Delta u = 2(\lambda u - f)$ and $\|\Delta u\|_p \leq 4\|f\|_p$. In particular for $1 < p < \infty$ there is a constant $C(p, d)$ such that for all $f \in L_p(R^d)$, $i, j = 1, \dots, d$ and $\lambda > 0$,

$$\|D_{x_i} D_{x_j} R_\lambda f\|_p \leq C(p, d) \|f\|_p$$

Let $\epsilon(p, d) = \frac{1}{d^2 C(p, d)}$. Then if $\sup_{i,j,x} |a_{i,j}(x) - \delta_{i,j}| \leq \epsilon(p, d)$

$$\begin{aligned} \left| (\lambda R_\lambda f)(x) - \frac{1}{2} \sum_{i,j} a_{i,j}(x) \frac{\partial^2 R_\lambda f}{\partial x_i \partial x_j}(x) - f(x) \right| &\leq \left| \frac{1}{2} \sum_{i,j} [(a_{i,j}(x) - \delta_{i,j})] \frac{\partial^2 R_\lambda f}{\partial x_i \partial x_j}(x) \right| \\ &\leq \frac{1}{2} \sup_{i,j,x} |a_{i,j}(x) - \delta_{i,j}| \sum_{i,j} \left| \frac{\partial^2 R_\lambda f}{\partial x_i \partial x_j}(x) \right| \end{aligned}$$

$$\left| (\lambda R_\lambda f)(\cdot) - \frac{1}{2} \sum_{i,j} a_{i,j}(\cdot) \frac{\partial^2 R_\lambda f}{\partial x_i \partial x_j}(\cdot) - f(\cdot) \right|_p \leq \frac{1}{2} \epsilon(p, d) d^2 C(d, p) \|f\|_p = \frac{1}{2} \|f\|_p$$

Denoting by

$$L = \frac{1}{2} \sum_{i,j} a_{i,j}(x) \frac{\partial^2 f}{\partial x_i \partial x_j}$$

and by

$$L_0 = \frac{1}{2} \Delta = \frac{1}{2} \sum_i \frac{\partial^2}{\partial x_i^2}$$

$$\|\lambda R_\lambda - LR_\lambda - I\|_{p \rightarrow p} \leq \frac{1}{2}$$

Implies $\lambda R_\lambda - LR_\lambda$ is an invertible operator $L_p \rightarrow L_p$. In particular since R_λ maps smooth functions into smooth functions, for any $\lambda > 0$ the range of $\lambda u - Lu$ as u varies over smooth functions with compact support is dense in L_p .

If P_1, P_2 are two solutions on $(C[[0, T]; R^d, \mathcal{F}_t)$ with $P_i[x(0) - x_0] = 1$ that make

$$X_u(t) = u(x(t)) - u(x(0)) - \int_0^t (Lu)(x(s))ds$$

a martingale for smooth u , then for smooth u

$$e^{-\lambda t}u(x(t)) - u(x(0)) - \int_0^t e^{-\lambda s}(\lambda u - Lu)(x(s))ds$$

is a martingale and

$$E^{P_i}[\int_0^\infty e^{-\lambda s}(\lambda u - Lu)(x(s))ds] = u(x_0)$$

We want to conclude that

$$E^{P_1}[\int_0^\infty e^{-\lambda s}f(x(s))ds] = E^{P_2}[\int_0^\infty e^{-\lambda s}f(x(s))ds]$$

for sufficiently many f and therefore by the uniqueness result for Laplace transforms conclude that

$$E^{P_1}[f(x(t))] = E^{P_2}[f(x(t))]$$

Since the set of functions f in the range of $\lambda I - L$ is dense in L_p we need to show that for any solution P the probability measure

$$\mu(A) = \int_0^\infty e^{-\lambda t}P[x(t) \in A]dt$$

is in L_q where $\frac{1}{p} + \frac{1}{q} = 1$. For Brownian motion the singularity at the origin is $\simeq |x|^{2-d}$. For $\lambda > 0$ there is exponential decay at ∞ so that $\mu \in L_q$ if $q(d-2) < d$ or $q < \frac{d}{d-2}$ or $p > \frac{d}{2}$. Fix $p_0 > \frac{d}{2}$. If P is a solution for

$$L = \frac{1}{2} \sum_{i,j} a_{i,j}(x) \frac{\partial^2 f}{\partial x_i \partial x_j}$$

and

$$|a_{i,j}(x) - \delta_{i,j}| \leq \epsilon(p_0, d)$$

so that for any $\lambda > 0$, the range of $\lambda u - Lu$ as u ranges over smooth functions with compact support is dense in L_{p_0} . Let q_0 satisfy $p_0^{-1} + q_0^{-1} = 1$. Let P be a solution for L with $P[x(0) = x_0] = 1$. Then we know that if $f = \lambda u - Lu$,

$$u(x_0) = E^P \left[\int_0^\infty e^{-\lambda s} f(x(s)) ds \right] = \int f(x) \mu_\lambda(dx)$$

where $\mu_\lambda(A) = \int_0^\infty e^{-\lambda t} P[x(t) \in A] dt$.

$$\begin{aligned} [(\lambda I - L_0)^{-1} f](x_0) &= E^P \int_0^\infty e^{-\lambda t} [(\lambda I - L)(\lambda - L_0)^{-1} f](x(t)) dt \\ &= E^P \int_0^\infty e^{-\lambda t} [(\lambda I - L_0 + L_0 - L)(\lambda - L_0)^{-1} f](x(t)) dt \\ &= E^P \int_0^\infty e^{-\lambda t} [I - (L - L_0)(\lambda - L_0)^{-1}] f(x(t)) dt \\ &= \langle f, \mu_\lambda \rangle - \langle (L - L_0)(\lambda - L_0)^{-1} f, \mu_\lambda \rangle \end{aligned}$$

Let us take supremum over f with $\|f\|_p \leq 1$. Then

$$\|\mu_\lambda\|_q \leq c(q, d) + \epsilon(p, d) \|\mu_\lambda\|_q$$

Either $\|\mu_\lambda\|_q = \infty$ or $\|\mu_\lambda\|_q \leq (1 - \epsilon(p, d))^{-1} c(p, d)$. We have an a priori bound, but we need to approximate P by P_n satisfying the same bound and pass to the limit. We know that there is a BM such that with $\sigma\sigma^* = a$, $\sigma = a^{\frac{1}{2}}$.

$$x(t) = x_0 + \int_0^t \sigma(x(s)) \cdot d\beta(s)$$

we can approximate by

$$x^n(t) = x_0 + \int_0^t \sigma^n(s, \omega) \cdot d\beta(s)$$

where $a^n(s, \omega) = \sigma^n(s, \omega)\sigma^n(s, \omega)^*$ satisfies $d^2 \sup_{i,j,t,\omega} |a_{i,j}^n(t, \omega) - \delta_{i,j}| \leq \epsilon(p, d)$.

Let $\mu_\lambda^n(A) = \int_0^\infty e^{-\lambda t} P[x^n(t) \in A] dt$.

$$\begin{aligned} [(\lambda I - L_0)^{-1} f](x_0) &= E^P \int_0^\infty e^{-\lambda t} \left[(\lambda I - \frac{1}{2} \sum_{i,j} a_{i,j}^n(t, \omega) \frac{\partial^2}{\partial x_i \partial x_j}) (\lambda - L_0)^{-1} f \right] (x^n(t)) dt \\ &= E^P \int_0^\infty e^{-\lambda t} \left[(\lambda I - L_0 + L_0 - \frac{1}{2} \sum_{i,j} a_{i,j}^n(t, \omega) \frac{\partial^2}{\partial x_i \partial x_j}) (\lambda - L_0)^{-1} f \right] (x^n(t)) dt \\ &= E^P \int_0^\infty e^{-\lambda t} \left[I - \frac{1}{2} \sum_{i,j} (a_{i,j}^n(t, \omega) - \delta_{i,j}) \frac{\partial^2}{\partial x_i \partial x_j} \right] (\lambda - L_0)^{-1} f(x^n(t)) dt \\ &= \langle f, \mu_\lambda^n \rangle - \langle \frac{1}{2} \sum_{i,j} (a_{i,j}^n(t, \omega) - \delta_{i,j}) \frac{\partial^2}{\partial x_i \partial x_j} (\lambda - L_0)^{-1} f, \mu_\lambda^n \rangle \end{aligned}$$

Let us again take supremum over f with $\|f\|_p \leq 1$. Then

$$\|\mu_\lambda^n\|_q \leq c(q, d) + \epsilon(p, d)\|\mu_\lambda^n\|_q$$

But $\|\mu_\lambda\|_q < \infty$ and $\|\mu_\lambda\|_q \leq (1 - \epsilon(p, d))^{-1}c(p, d)$ We can now pass to the limit.

A similar argument holds for time dependent situation where we have $a_{i,j}(t, x)$ and then the Laplace transform is not useful. We need to solve the Cauchy problem

$$u_t + \frac{1}{2} \sum_{i,j} a_{i,j}(t, x) \frac{\partial^2 u}{\partial x_i \partial x_j} = -f(t, x); \text{ for } s < T, \text{ and } u(T, x) = 0$$

Then if P is a solution with $P[x(s) = x_0] = 1$,

$$u(t, x(t)) - u(s, x(s)) - \int_s^t u_\sigma(\sigma, x(\sigma))d\sigma - \int_s^t \frac{1}{2} \sum_{i,j} a_{i,j}(\sigma, x(\sigma)) \frac{\partial^2 u}{\partial x_i \partial x_j}(\sigma, x(\sigma))d\sigma$$

i.e.

$$u(t, x(t)) - u(s, x(s)) + \int_s^t f(\sigma, x(\sigma))d\sigma$$

is a martingale with respect to any solution P with $P[x(s) = x_0] = 1$. Equating expectations at $t = s$ and $t = T$

$$u(s, x_0) = E^P \left[\int_s^T f(\sigma, x(\sigma))d\sigma \right]$$

If enough expectations are determined, then P is unique. The equation

$$u_t + \frac{1}{2} \Delta u = -f; u(T, x) = 0$$

has the solution

$$u(s, x) = \int_s^T \int_{R^d} f(t, y) \frac{1}{(2\pi(t-s))^{\frac{d}{2}}} \exp\left[-\frac{(y-x)^2}{2(t-s)}\right] dy dt$$

One has analogs of Calderon-Zygmund estimates (Ben Franklin Jones) that estimate $\|u_t\|_p, \|u_{x_i, x_j}\|_p$ on $R^d \times [0, T]$ in terms of $\|f\|_p$ on $R^d \times [0, T]$ for $1 < p < \infty$. This allows the perturbation to work. We need to pick a p such that for $q = \frac{p}{p-1}$

$$\int_0^T \int_{R^d} \frac{1}{(2\pi t)^{\frac{dq}{2}}} \exp\left[-\frac{qy^2}{2t}\right] dy dt < \infty$$

i.e $d(q-1) < 2$ or $q < 1 + \frac{2}{d}$ or $p > \frac{1+\frac{2}{d}}{\frac{2}{d}} = \frac{d+2}{2}$